Aspects of LSS Polarized DIS analysis

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with thanks to

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- LSS polarized PDFs and comparison with DSSV and AAC

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$$\frac{g_1}{F_1} = \frac{A_{\parallel}}{d} + \frac{2Mxg_2}{(E + E'\cos\theta)F_1}$$

$$\approx \frac{A_{\parallel}}{d}$$

$$A_1 = \frac{A_{\parallel}}{D} - \eta A_2$$

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$$A_1 = \frac{g_1 - \gamma^2 g_2}{F_1}$$

and ignore g_2 term, or

replace g_2 in terms of A_1 and A_2

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$$\frac{g_1}{F_1} \approx A_1 \quad \text{or} \quad \frac{A_1}{1 + \gamma^2}$$

$$\gamma^2 = \frac{4M^2x^2}{Q^2}$$

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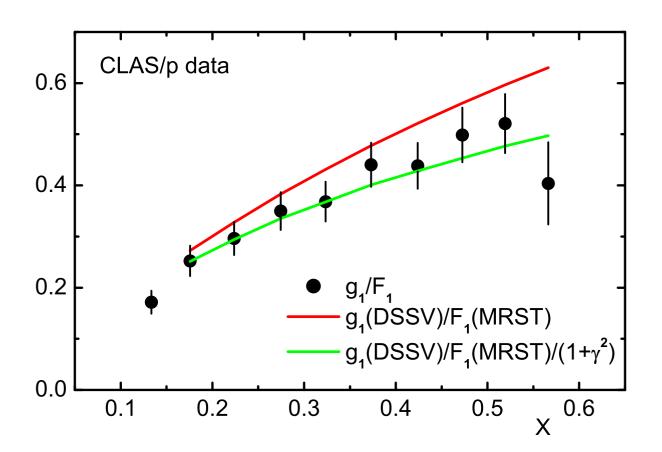
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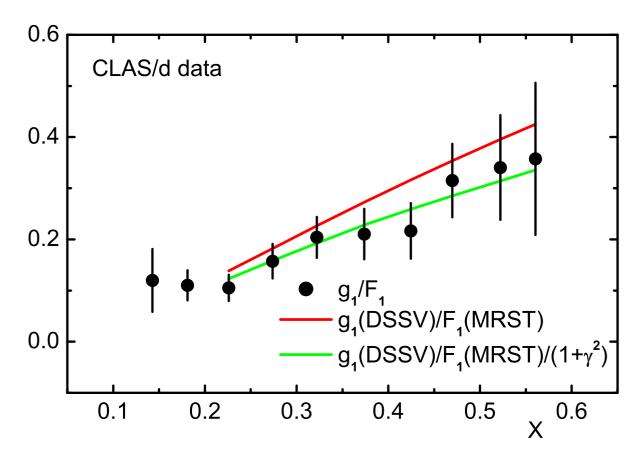
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Effect on PDFs ??????? Look at CLAS proton and neutron data:

Compare
$$\left(\frac{g_1}{F_1}\right)_{DSSV}/(1+\gamma^2)$$
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Comparison of $\chi^2 s$ for fit to $\left(\frac{g_1}{F_1}\right)_{Expt}$

Expt	$\left(\frac{g_1}{F_1}\right)_{DSSV}/(1+\gamma^2)$	$\left(\frac{g_1}{F_1}\right)_{DSSV}$
р	5.9	20
n	2.5	8.2

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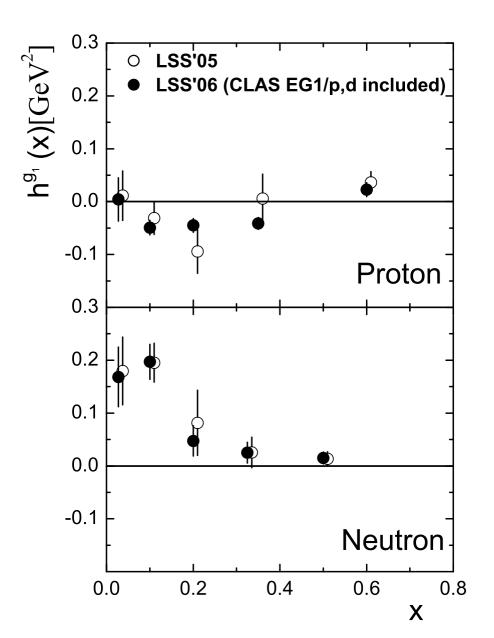
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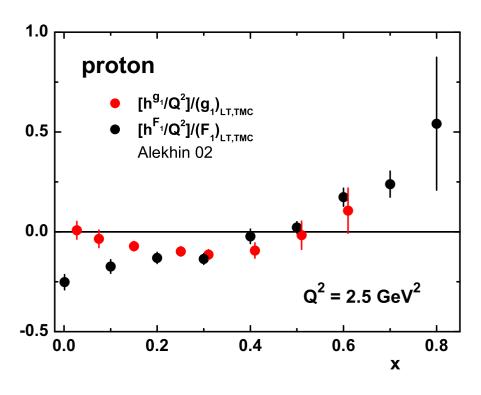
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$$\left[\frac{g_1}{F_1}\right]^{EXP} \approx \frac{g_1^{LT}}{F_1^{LT}} \left[1 + \frac{g_1^{HT}}{g_1^{LT}} - \frac{F_1^{HT}}{F_1^{LT}}\right] \approx \frac{g_1^{LT}}{F_1^{LT}}$$

provided there is a cancellation between $\frac{g_1^{HT}}{g_1^{LT}}$ and $\frac{F_1^{HT}}{F_1^{LT}}$.





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How will this affect the DSSV PDFs??

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$$\Delta q(x)|_B = T_{B \leftarrow A} \, \Delta q(x)|_A. \tag{4}$$

Suppose now that $T_{B\leftarrow A}$ is known to NLO accuracy, and the parton densities are extracted from the data, *independently*, in NLO, using schemes A and B, with results $\Delta q(x)|_{A,B}^{data}$, respectively.

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Thus the ratio

$$\frac{\Delta q(x)|_B^{data} - T_{B \leftarrow A} \Delta q(x)|_A^{data}}{\Delta q(x)|_B^{data} + T_{B \leftarrow A} \Delta q(x)|_A^{data}}$$

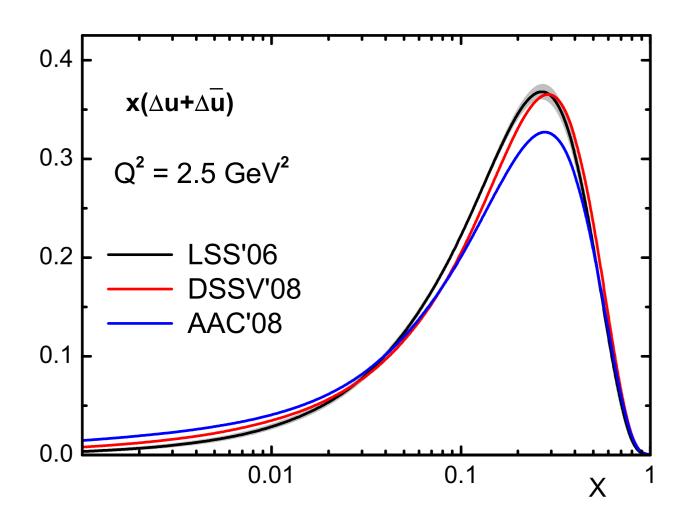
gives some indication of the reliability of the parton densities.

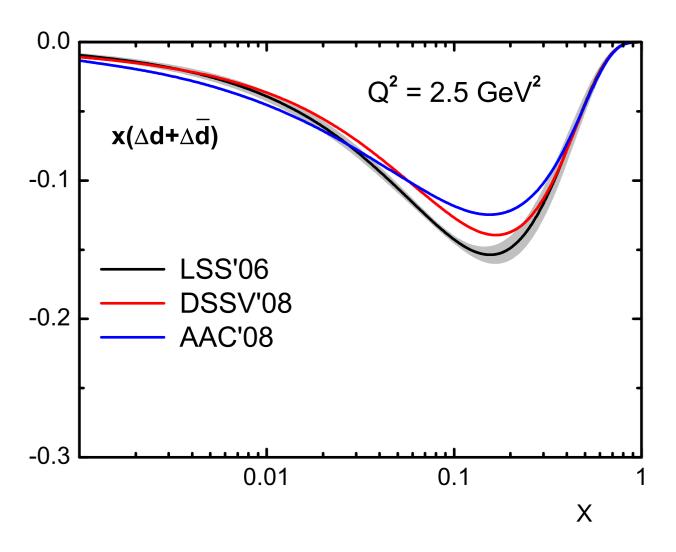
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It was shown that a positive value for the first moment would imply a huge breaking of $SU(3)_F$ invariance, far greater than the $\pm 10\%$ breaking estimated by Ratcliffe

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However, the DSSV combined analysis (DIS, SIDIS, $pp \to \pi$) also finds positive values for $\Delta s(x) + \Delta \bar{s}(x)$ for $x \geq 0.03$, yet ends up with a negative first moment $\Delta S = -0.114$ at $Q^2 = 10 \, GeV^2$.

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COMPASS (Windmolders) study dependence of $\Delta s(x) + \Delta \bar{s}(x)$ on the choice of fragmentation functions.

LO SIDIS K^+ and K^- production: $0.004 < x \le 0.3$

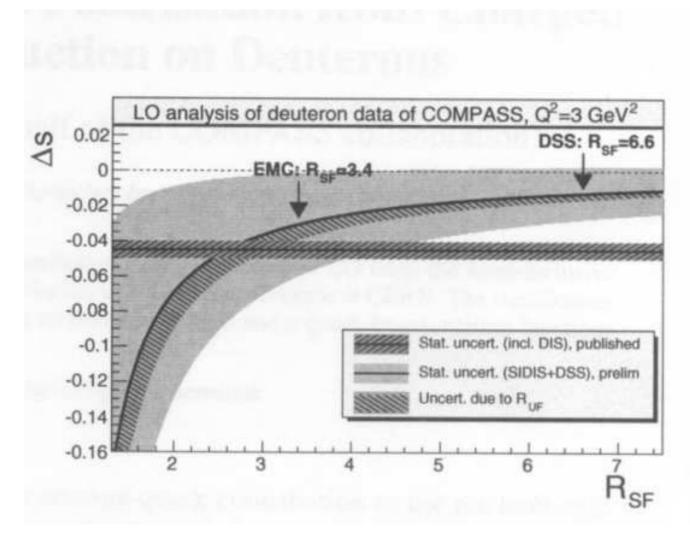
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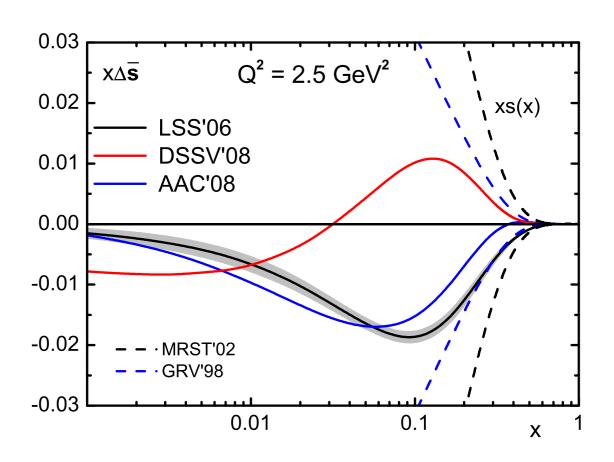
Plot integral over measured range vs $R_{S/F}$ Result sensitive to $R_{S/F}$



Main curve uses $R_{U/F}=0.14$ (DSS value); hatched uses 0.35 (SMC value).

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COMPASS
$$\Delta\Sigma(Q^2 = 3) = 0.35 \pm 0.06$$

HERMES $\Delta\Sigma(Q^2 = 5) = 0.33 \pm 0.04$

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Two standard deviations difference! No explanation.

The polarized gluon density

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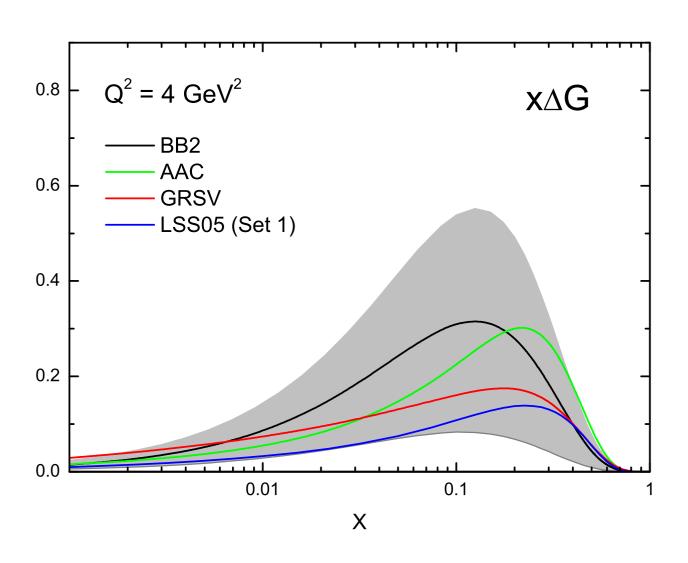
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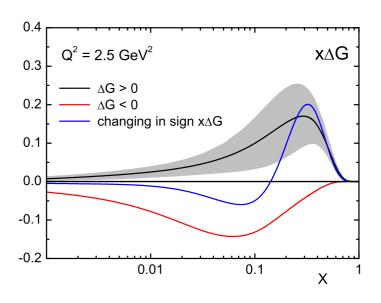
For a long time all analyses seemed to indicate that $\Delta G(x)$ was a positive function of x.

ΔG a few years ago:



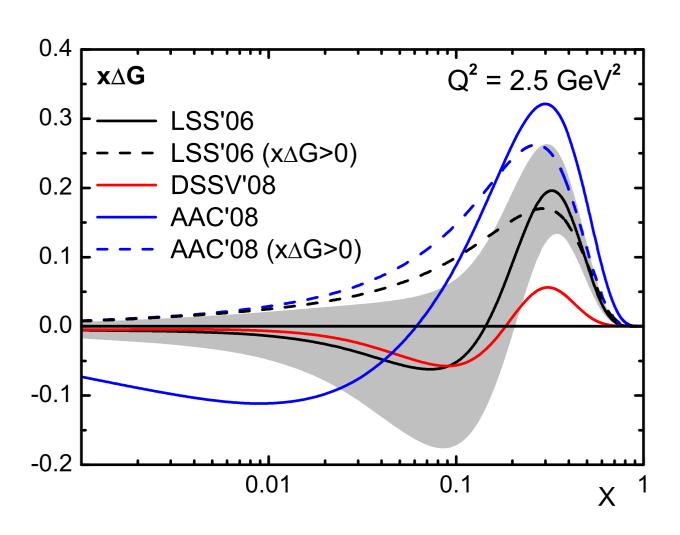
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In all fits, and irrespective of the form of the gluon density, the magnitude is always found to be very small. In all fits, and irrespective of the form of the gluon density, the magnitude is always found to be very small. Typically one has $|\Delta G|\approx 0.29\pm 0.32$,

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- Need to understand disagreements in first moment $\Delta\Sigma$ obtained from HT expansions.

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- (2) IF assume $\Delta\Sigma|_{JET} \Leftrightarrow 2S_z^{quarks} \approx 60\%$ then need $\Delta G \approx 1.7$ at $Q^2 = 1 GeV^2$

Much bigger than present values!